

chloride (1.75 Å)¹² and markedly longer than in diboron tetrafluoride (1.67 Å).¹³ The B-P and P-F bond distances are very similar to those observed in trifluorophosphine-borane (F₃P·BH₃),¹⁴ while the P-F distances are equal and the B-P distance is only slightly shorter in the borane adduct (0.012 Å). However, the trifluorophosphine group in the borane adduct has C_{3v} symmetry but not in the refined structure of model A. A planar projection of the trifluorophosphine group shows two FPF angles of 122 ± 0.6°, which reduces the symmetry to C_{2v}. The deviation from C_{3v} symmetry is such as to increase the nearest fluorine-hydrogen nonbonded distance.

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Model A does show one unusual structural feature. The HBH angle of 146 ± 12° is much larger than that reported for diborane(6) and tetraborane(10), 122°,⁹ in trifluorophosphine-borane and carbonyl-borane, these are 115 and 113°,¹⁵ respectively. The small BBH angle of 92.3° is associated with a short nonbonded boron-hydrogen distance of 2.23 Å. In contrast, the hydrogen positions in model B are not significantly different from those reported for diborane(6).

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The Crystal Structure of Tetraethylammonium Hexabromoantimonate(V), (C₂H₅)₄NSbBr₆¹

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Tetraethylammonium hexabromoantimonate(V) has been prepared and its crystal structure determined by three-dimensional X-ray analysis. The dark red-brown compound crystallizes in the tetragonal space group *I*4₁*md* with four molecules in a unit cell of dimensions *a* = *b* = 8.7008 (7) and *c* = 24.797 (3) Å. Data were collected with an automated diffractometer using the θ -2 θ step scan technique and Mo K α radiation (λ 0.7107 Å). The structure was solved by Patterson and Fourier methods and refined by full-matrix least-squares techniques to a discrepancy index *R* = 0.042 for the 276 independent reflections used in the analysis. The structure consists of slightly distorted octahedral Sb^VBr₆⁻ anions and (C₂H₅)₄N⁺ cations which are twofold disordered about intersecting mirror planes. The average Sb-Br distances corrected for thermal motion assuming a riding model are 2.561 (5), 2.564 (5), and 2.570 (4) Å. The only close Br-Br van der Waals contact is along the *a* axis at 3.584 Å. The similar sizes of the tetraethylammonium and the more spherical hexabromoantimonate(V) ions permit efficient packing which contributes to the stability of this compound.

Introduction

The crystal structure investigation of tetraethylammonium hexabromoantimonate(V) was undertaken as part of a series of investigations of halo-coordinated antimony compounds being performed in this laboratory.⁴⁻¹¹ A series of intervalence antimony bromide complexes of varying stoichiometries has been investigated to relate, if possible, their crystal structures and charge transfer properties and to evaluate the effects of cation size, type, and stereochemistry on the antimony bromide anion and on the resulting solid-state structure. The structure of the tetraethylammonium salt was under-

taken as a result of preliminary investigations⁸ on the physical and chemical properties of this material and other intensely colored R₄Sb₃Br₆ complexes. The tetraethylammonium complex was shown to be unusually stable in comparison to the other antimony bromide complexes studied, being stable in air and more slowly hydrolyzed in water. Although the tetraethylammonium complex is also deeply colored,¹² appearing a dark red-brown, the preliminary X-ray results indicated the stoichiometry was RSb^VBr₆ with only four formula units per unit cell. Intervalence charge transfer could not occur if indeed all the SbBr₆ species present were equivalent. Also, the saturated cation would not be expected to participate in any other possible type of charge-transfer path.

Experimental Section

Tetraethylammonium hexabromoantimonate(V), (C₂H₅)₄NSbBr₆, mol wt 479.5, crystallizes in the tetragonal space group *I*4₁*md*, with *F*(000) = 1344 e, *Z* = 4, *a* = *b* = 8.7008 (7) Å, *c* = 24.797 (3) Å, temperature ~24°, *V* = 1877.3 Å³, *d*_m = 2.62 g/cm³, *d*_o = 2.59 g/cm³, Mo K α (λ 0.7107 Å), and μ = 150.7 cm⁻¹.

The (C₂H₅)₄NSbBr₆ salt was prepared by the method of Law-

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TABLE I
FINAL ATOMIC POSITIONAL^a AND THERMAL PARAMETERS^b AND STANDARD ERRORS^c FOR SbBr_6^-

Atom	Position	<i>x</i>	<i>y</i>	<i>z</i>	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Sb	4a	0 ^d	0 ^d	0 ^d	87 (5)	117 (6)	13 (1)	0 ^d	0 ^d	0 ^d
Br(1)	8b	0 ^d	0.2034 (7)	0.0732 (2)	244 (12)	217 (11)	22 (1)	0 ^d	0 ^d	-29 (3)
Br(2)	8b	0 ^d	-0.2057 (6)	-0.0734 (2)	135 (8)	155 (8)	20 (1)	0 ^d	0 ^d	-16 (2)
Br(3)	8b	0.2940 (4)	0 ^d	-0.0027 (3)	97 (6)	175 (7)	24 (1)	0 ^d	-13 (3)	0 ^d

^a Positional parameters are in fractional unit cell coordinates. ^b β 's $\times 10^4$; the form of the anisotropic temperature factor is $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. ^c Estimated standard deviations are given in parentheses for the least significant figures. ^d Atomic parameter fixed by symmetry.

ton.⁶ Microscopic examination revealed the fourfold symmetry of the crystals, square pyramidal with sharply defined faces. Although the tetraethylammonium salt was observed to be air stable, crystals were selected and mounted in thin-walled Lindemann glass capillaries to limit exposure in the atmosphere. Preliminary Weissenberg and precession photographs exhibited $4/mmm$ Laue symmetry, indicating a tetragonal space group. The following systematic absences were observed: hkl when $h + k + l = 2n + 1$ and hkl when $2h + l = 4n + 1$. The absences are consistent with either space group $I4_1md$ or $I\bar{4}2d$. The unit cell parameters and their standard deviations were obtained by a least-squares fit to 12 independent reflection angles whose centers were determined by a left-right, top-bottom beam-splitting technique using a previously aligned Hilger-Watts four-circle diffractometer (Mo $K\alpha$ radiation, λ 0.71069 Å). The observed density was determined by flotation techniques using ethylene bromide and 1,1,2,2-tetrabromoethane solutions.

A crystal having approximate dimensions $0.10 \times 0.10 \times 0.09$ mm was mounted on a glass fiber such that the *c* axis (0.09 mm) was nearly coincident with the spindle axis. Data were collected at room temperature utilizing a Hilger-Watts four-circle diffractometer equipped with scintillation counter and using Zr-filtered Mo $K\alpha$ radiation. Within a 2θ sphere of 60° all data in one full octant were recorded using a θ - 2θ step scan technique with a takeoff angle of 4.5° . Symmetric scan ranges of 1.00° in 2θ at low 2θ values to 2.20° at the high 2θ limit were used. Stationary-crystal, stationary-counter background measurements were made at the beginning and end of the step scan, each measurement being made for half the total scan time. The counting rate used was 0.2048 sec/step of 0.02° in 2θ . The rest of the experimental arrangement used has already been discussed in some detail.¹³ A total of 1535 reflections were measured in this way. Based on the measurements of three standard reflections, no decomposition correction was made; joint decomposition was less than 3%.

The intensity data were also corrected for Lorentz-polarization effects and for effects due to absorption ($\mu = 150.7 \text{ cm}^{-1}$). The absorption correction¹⁴ was made using ABCOR¹⁵ the maximum and minimum transmission factors were 0.414 and 0.322, respectively. The error assigned each intensity measurement was calculated from the estimated statistical and nonstatistical errors by

$$[\sigma(I)]^2 = [C_T + C_B + (0.03C_T)^2 + (0.05C_B)^2 + (0.04C_R)^2]/A^2$$

where C_T , C_B , C_R , and A are the total count, background count, net count, and transmission factor, respectively. The equivalent values of F_o^2 were then averaged. The estimated standard deviation in each structure factor was calculated from the mean deviation of intensity by the method of finite differences.¹⁶ The reciprocals of the structure factor variances were used as weights in the least-squares refinement. Based on the measurements of symmetry extinct data, it was decided that only those reflections with $I \geq 2\sigma(I)$ would be considered observed. The results reported are based on the remaining 276 independent reflections.

Solution and Refinement of the Structure

With only four formula units per unit cell, the anti-

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(15) In addition to local programs written at this laboratory, the authors wish to acknowledge the use of ABCOR (Busing and Levy), ORTEP (C. K. Johnson), ORFFE and a local modification of ORFLS (Busing, Martin, and Levy), and LCR2 (D. E. Williams).

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TABLE II
FINAL ATOMIC PARAMETERS AND STANDARD ERRORS FOR $(\text{C}_2\text{H}_5)_4\text{N}^+$

Atom	Position	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> , Å ²
N	4a	0 ^b	0 ^b	0.4333 (16)	4.0 (9)
C(1)	8b	0.2158 (52)	0 ^b	0.3666 (17)	5.9 (10)
C(2)	8b	0.2170 (47)	0 ^b	0.5064 (23)	5.5 (10)
C(3) ^a	16c	0.1005 (64)	0.1185 (62)	0.3976 (18)	4.1 (11)
C(4)	16c	0.0999 (63)	-0.0998 (70)	0.4684 (21)	5.0 (11)

^a Atoms C(3) and C(4) are disordered and refined with half-occupancy. ^b Atomic parameter fixed by symmetry.

TABLE III
OBSERVED AND CALCULATED STRUCTURE FACTORS ($\times 10$)

L	h	k	l	Observed	Calculated	L	h	k	l	Observed	Calculated			
0	0	0	0	3127	3028	2	0	2	2	1234	1233			
0	0	0	1	1724	1675	4	0	0	505	505	6	6	671	581
0	0	1	1	2033	1975	3	1	1	555	586	3	1	555	586
0	1	0	0	1042	1035	2	2	2	1751	1853	4	2	1050	1074
0	1	0	1	510	565	4	2	0	925	864	1	0	2444	2575
0	1	1	0	605	594	6	2	0	955	864	3	0	2243	2256
0	1	1	1	444	395	4	4	0	925	567	5	0	1935	1971
1	0	0	2	2987	2923	7	1	0	520	597	7	0	1074	1074
1	0	1	0	1291	1261	2	2	0	801	851	6	4	807	450
1	0	1	1	1286	1305	4	2	1	1116	1107	8	8	584	199
1	0	2	0	4015	3905	6	2	0	564	535	4	4	1006	980
1	0	2	1	2453	2488	3	3	0	554	542	6	1	850	844
1	0	2	2	564	555	4	4	0	1361	1335	3	2	864	849
1	0	3	0	2649	2728	6	4	0	779	768	5	2	974	998
1	0	3	1	1366	1395	3	0	1372	1409	4	2	972	1005	
1	0	3	2	469	459	5	0	1842	1759	6	3	705	669	
1	0	3	3	1157	1217	7	0	1063	1084	5	4	1059	1013	
1	0	4	0	809	711	1	0	1824	1905	2	1	1338	1331	
1	0	4	1	0	0	2	0	2480	2457	4	1	711	758	
1	0	4	2	0	0	3	0	619	753	6	1	534	564	
1	0	4	3	0	0	3	0	943	990	3	2	1933	1877	
1	0	4	4	0	0	2	1	1391	1410	5	2	472	486	
1	0	4	5	0	0	1	1	643	648	7	2	478	1050	
1	0	4	6	0	0	6	1	3648	3710	4	3	1170	1115	
1	0	4	7	0	0	3	2	2285	2227	6	3	984	960	
1	0	4	8	0	0	5	2	736	682	8	3	547	454	
1	0	4	9	0	0	7	2	542	529	5	4	562	474	
1	0	4	10	0	0	2	2	1095	1120	5	4	518	607	
1	0	4	11	0	0	4	3	1569	1582	7	6	418	607	
1	0	4	12	0	0	6	3	935	957	5	5	856	838	
1	0	4	13	0	0	7	4	610	725	7	5	524	391	
1	0	4	14	0	0	2	0	1861	1894	6	0	597	581	
1	0	4	15	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	16	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	17	0	0	3	1	984	960	3	1	797	809	
1	0	4	18	0	0	8	3	547	454	5	1	1308	1330	
1	0	4	19	0	0	5	4	562	474	2	2	471	408	
1	0	4	20	0	0	6	4	518	607	5	3	700	679	
1	0	4	21	0	0	5	5	856	838	4	3	605	592	
1	0	4	22	0	0	7	5	524	391	7	5	524	391	
1	0	4	23	0	0	2	0	1861	1894	6	0	597	581	
1	0	4	24	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	25	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	26	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	27	0	0	7	0	694	799	7	0	694	799	
1	0	4	28	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	29	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	30	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	31	0	0	7	0	694	799	7	0	694	799	
1	0	4	32	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	33	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	34	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	35	0	0	7	0	694	799	7	0	694	799	
1	0	4	36	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	37	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	38	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	39	0	0	7	0	694	799	7	0	694	799	
1	0	4	40	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	41	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	42	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	43	0	0	7	0	694	799	7	0	694	799	
1	0	4	44	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	45	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	46	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	47	0	0	7	0	694	799	7	0	694	799	
1	0	4	48	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	49	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	50	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	51	0	0	7	0	694	799	7	0	694	799	
1	0	4	52	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	53	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	54	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	55	0	0	7	0	694	799	7	0	694	799	
1	0	4	56	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	57	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	58	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	59	0	0	7	0	694	799	7	0	694	799	
1	0	4	60	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	61	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	62	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	63	0	0	7	0	694	799	7	0	694	799	
1	0	4	64	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	65	0	0	3	1	1170	1115	3	0	1460	1422	
1	0	4	66	0	0	3	1	984	960	3	0	1002	1019	
1	0	4	67	0	0	7	0	694	799	7	0	694	799	
1	0	4	68	0	0	1	1	918	931	1	0	1595	1675	
1	0	4	69	0	0	3	1	1170	1115	3	0	1460	1422	
1														

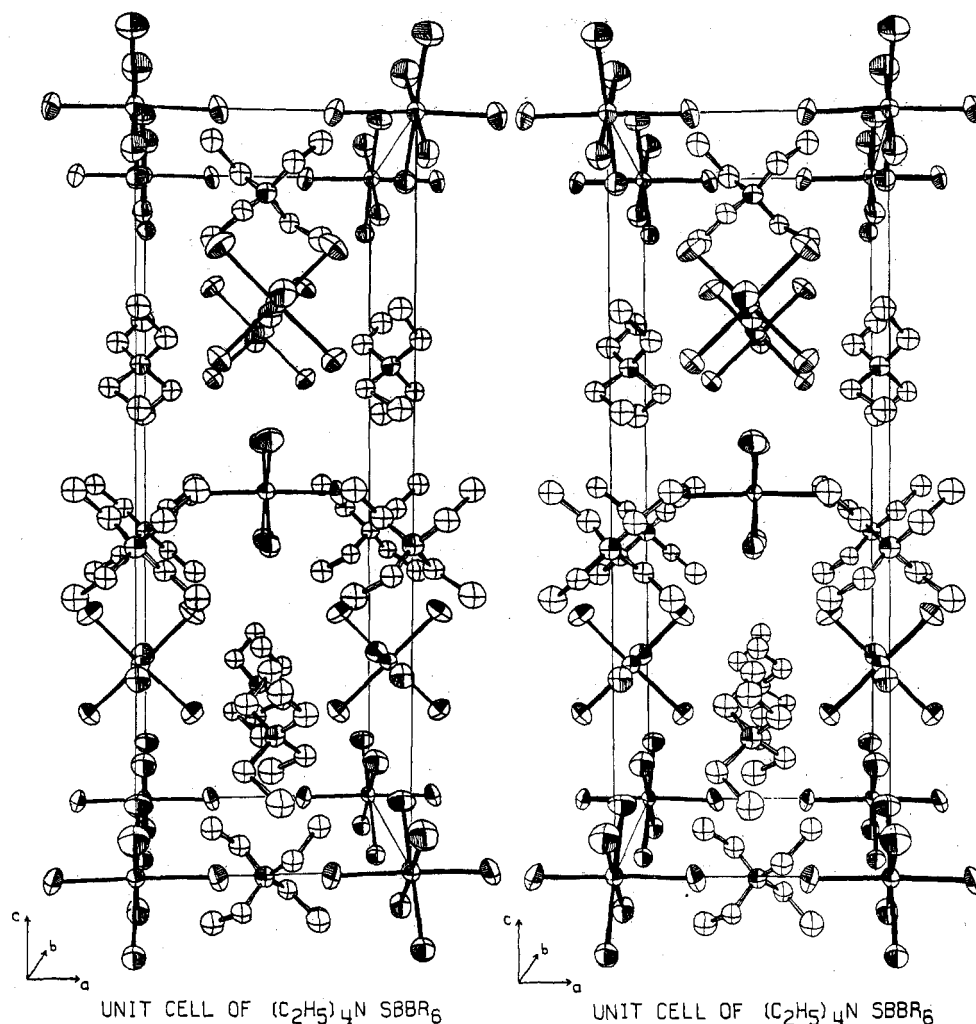


Figure 1.—Stereogram of unit cell showing packing of $(\text{C}_2\text{H}_5)_4\text{NSbBr}_6$ units; hydrogens not shown.

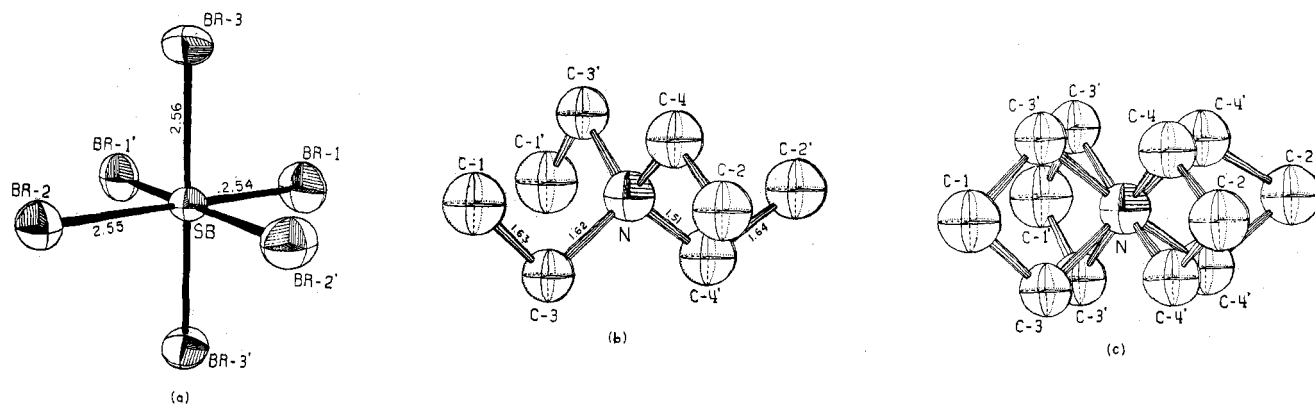


Figure 2.—Ions present in $(\text{C}_2\text{H}_5)_4\text{NSbBr}_6$: (a) SbBr_6^- ; (b) $(\text{C}_2\text{H}_5)_4\text{N}^+$; (c) disordered $(\text{C}_2\text{H}_5)_4\text{N}^+$.

A tetraethylammonium group has two preferred orientations, the swastika configuration or a trans arrangement. Although found in a trans arrangement, the inner carbon atoms do not lie on the mirror planes and are twofold disordered. These inner carbons were refined with half-occupancy in a general 16-fold position. The outer carbons appeared to lie on the mirror planes and were treated as ordered and restrained to these mirror planes. Subsequent refinement using anisotropic thermal parameters for only the heavy atoms lowered the R factor to 0.042 and $R_w = 0.044$.

The relativistic Hartree-Fock X-ray scattering factors for neutral atoms of Doyle and Turner¹⁷ were used with those of antimony and bromine modified for the real and imaginary parts of anomalous dispersion.¹⁸ Based on the agreement of the large structure factors, no extinction correction was necessary. A final electron density difference map showed no peak heights

(17) P. A. Doyle and P. S. Turner, *Acta Crystallogr., Sect. A*, **24**, 390 (1968).

(18) D. H. Templeton, "International Tables for X-Ray Crystallography," Vol. III, Kynoch Press, Birmingham, England, 1962, pp 215, 216, Table 3.3.2C.

greater than $0.3 \text{ e}^-/\text{\AA}^3$. The final standard deviation for an observation of unit weight $([\sum w\Delta^2/(NO - NV)]^{1/2})$ where $\Delta = |F_o| - |F_c|$, NO is the number of observations (276), and NV is the number of variables (38) was 0.95 electron. During the final cycle, the largest shift in any parameter was less than 0.01 times its own σ . The final positional and thermal parameters are given in Tables I and II, along with their standard deviations as derived from the inverse matrix of the final least-squares cycle. In Table III are listed the magnitudes of the observed and calculated structure factors in electrons $\times 10$.

Discussion

The crystal structure of tetraethylammonium hexabromoantimonate(V) is shown in Figure 1. Bond distances and angles of interest are given in Table IV and

TABLE IV
SELECTED INTERATOMIC BOND DISTANCES
AND ANGLES FOR $(\text{C}_2\text{H}_5)_4\text{NSbBr}_6^a$

Atoms	Distance, \AA	Atoms	Angle, deg
Sb-Br(1)	2.536 (5)	Br(1)-Sb-Br(1)'	88.5 (3)
Sb-Br(2)	2.553 (5)	Br(1)-Sb-Br(2)'	91.2 (1)
Sb-Br(3)	2.559 (4)	Br(1)-Sb-Br(3)	91.1 (1)
Br(1)-Br(1)'	3.540 (12)	Br(2)-Sb-Br(2)'	89.0 (3)
Br(1)-Br(2)'	3.636 (5)	Br(2)-Sb-Br(3)	88.9 (1)
Br(1)-Br(3)	3.636 (5)	Br(1)-Sb-Br(2)	179.8 (5)
Br(2)-Br(2)'	3.579 (10)	Br(3)-Sb-Br(3)'	177.0 (3)
Br(2)-Br(3)	3.581 (5)	C(3)-N-C(4)	112 (3)
Br(3)-Br(3)' _i	3.584 (5)	C(3)-N-C(4)'	105 (3)
Br(1)-Br(2)' _{ii}	4.053 (6)	C(3)-N-C(3)'	113 (4)
N-C(3)	1.62 (5)	C(3)'-N-C(4)	105 (3)
N-C(4)	1.51 (5)	C(4)-N-C(4)'	109 (5)
C(3)-C(1)	1.63 (6)	N-C(3)-C(1)	101 (3)
C(4)-C(2)	1.64 (7)	N-C(4)-C(2)	113 (4)
Sb-Br(1)	2.561 (5) ^b		
Sb-Br(2)	2.564 (5) ^b		
Sb-Br(3)	2.570 (4) ^b		

^a Primed atoms refer to the symmetry-related atom in the group (Figure 2). Other symmetry operations referred to: (i) $1 + x, y, z$; (ii) $x, 1/2 + y, 1/4 + z$. ^b Interatomic distance corrected for thermal motion using a riding model where the second atom is assumed to ride on the first.

Figure 2. The $\text{Sb}^{\text{V}}\text{Br}_6^-$ ion has crystallographic C_{2v} symmetry but is somewhat distorted from O_h symmetry.

The most significant deviation involves the Br(3)-Sb-Br(3)' angle which is $177.0 (3)^\circ$. This slight distortion can be ascribed to packing effects since the closest approach between anions is $3.584 (5) \text{ \AA}$ along the a direction (Br(3)-Br(3)'_i) which is significantly shorter than the 3.9-\AA sum of the van der Waals radii.¹⁹ The average Sb-Br bond length is $2.549 (5) \text{ \AA}$ before correction for thermal motion and $2.565 (5) \text{ \AA}$ when corrected assuming a riding model. These averages are in good agreement with those previously reported.⁹ The tetraethylammonium ion has the trans configuration in which the ethyl groups lie on intersecting mirror planes (C_{2v} symmetry) as required by this space group. However, the inner carbon atoms do not lie on the mirror planes and are therefore disordered with apparent D_{2h} symmetry, as shown in Figure 2. Disorder within the swastika configuration of the tetraethylammonium ion has also been reported.²⁰ The long bond lengths indicate that the light-atom positions are not well defined and reflect both the disorder and the heavy-atom nature of this problem.

The crystal structure (Figure 1) can be viewed as an efficient packing arrangement of the rather spherical hexabromoantimonate(V) ions and of the nearly equal in size but slightly flattened tetraethylammonium ions. The similar sizes of these two large, rather diffuse ions contribute to the crystal stability. The usual type of intervalence charge transfer cannot occur in this structure. The structure consists of only one kind of SbBr_6^- species, has a saturated cation, and has only the one bromine---bromine contact (Br(3)-Br(3)'_i, 3.58 \AA) which is less than the sum of the van der Waals radii. This distance does not appear to be short enough for any type of interspecies charge-transfer interaction since the Br(3)-Br(2) and Br(3)-Br(1) intraion distances are 3.58 and 3.64 \AA , respectively. We believe, therefore, that the deep color of this complex probably results from normal charge transfer of the intraspecies ligand to metal type, involving transitions between molecular orbitals of the $\text{Sb}^{\text{V}}\text{Br}_6^-$ species.

(19) L. Pauling, "The Nature of the Chemical Bond," Cornell University Press, Ithaca, N. Y., 1960, p 260.

(20) G. D. Stucky, J. B. Folkers, and T. J. Kistenmacher, *Acta Crystallogr.*, **23**, 1064 (1967).

Notes

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Bromine Perchlorate

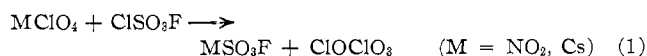
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Recently we reported the synthesis of a novel chlorine oxide, chlorine perchlorate.¹ This preparation was

(1) C. J. Schack and D. Pilipovich, *Inorg. Chem.*, **9**, 1387 (1970).

accomplished by the reaction



It has now been found that the related bromine compound bromine perchlorate can be prepared by this method using bromine(I) fluorosulfate.



In addition, a second method involving the oxidation of elemental bromine with chlorine perchlorate was discovered

